

Physics 606 Exam 1 Solution

These are the main steps.

Your solution is likely to be more detailed.

1. Use $x = \sqrt{\frac{\hbar}{2m\omega}} (a^\dagger + a)$, $p = i\sqrt{\frac{m\hbar\omega}{2}} (a^\dagger - a)$

$$(\Delta x)^2 = \langle n | x^2 | n \rangle - (\langle n | x | n \rangle)^2 \quad [(\Delta x)^2 \equiv \langle (x - \langle x \rangle)^2 \rangle]$$

$$(\Delta p)^2 = \langle n | p^2 | n \rangle - (\langle n | p | n \rangle)^2 = \langle x^2 - 2x\langle x \rangle + \langle x \rangle^2 \rangle = \langle x^2 \rangle - \langle x \rangle^2]$$

$$a |n\rangle = \sqrt{n} |n-1\rangle, \quad a^\dagger |n\rangle = \sqrt{n+1} |n+1\rangle$$

Then $\langle n | x | n \rangle = \sqrt{\frac{\hbar}{2m\omega}} (\underbrace{\sqrt{n+1} \langle n | n+1 \rangle}_{=0} + \underbrace{\sqrt{n} \langle n | n-1 \rangle}_{=0}) = 0$

and similarly $\langle n | p | n \rangle = 0$, so

(a) $(\Delta x)^2 = \langle x^2 \rangle = \frac{\hbar}{2m\omega} \langle n | (a^{\dagger 2} + a^\dagger a + a a^\dagger + a^2) | n \rangle$

$$= \frac{\hbar}{2m\omega} (0 + \sqrt{n} \cdot \sqrt{n} + \sqrt{n+1} \cdot \sqrt{n+1}) \underbrace{\langle n | n \rangle}_{=1}$$

$$= \frac{\hbar}{2m\omega} (2n+1)$$

(b) Similarly, $(\Delta p)^2 = \frac{m\hbar\omega}{2} (2n+1)$

Then $|\Delta p \Delta x| = \sqrt{\frac{m\hbar\omega}{2}} \sqrt{\frac{\hbar}{2m\omega}} \sqrt{2n+1} \sqrt{2n+1}$

$$= \hbar (n + \frac{1}{2})$$

2. Fourier transform of $f(x)$:

$$f(k) = \int dx f(x) e^{-ikx}$$

$$\rightarrow f(\vec{k}) = \int d^3x f(\vec{r}) e^{-i\vec{k} \cdot \vec{r}}$$

(a) so, for $\delta^{(3)}(\vec{r} - \vec{r}_0)$, $f(\vec{k}) = \int d^3x \delta^{(3)}(\vec{r} - \vec{r}_0) e^{-i\vec{k} \cdot \vec{r}}$

$$= e^{-i\vec{k} \cdot \vec{r}_0}$$

(b) With this convention,

$$f(x) = \int \frac{dk}{2\pi} f(k) e^{ikx}$$

$$f(\vec{r}) = \int \frac{d^3k}{(2\pi)^3} f(k) e^{i\vec{k} \cdot \vec{r}}$$

$$\delta(\vec{r} - \vec{r}_0) = \int \frac{d^3k}{(2\pi)^3} e^{i\vec{k} \cdot (\vec{r} - \vec{r}_0)}$$

$$3. (a) \boxed{[X_H(t), P_H(t)]} = [x(t), p(t)] \cos^2 \omega t - [p(t), x(t)] \sin^2 \omega t$$

$$= i\hbar (\cos^2 \omega t + \sin^2 \omega t) = \boxed{i\hbar}$$

$$(b) \frac{dX_H(t)}{dt} = \frac{1}{i\hbar} [X(t), H(t)] + \frac{\partial X_H(t)}{\partial t}, \quad H(t) = \frac{p(t)^2}{2m} + \frac{1}{2} m \omega^2 x(t)^2,$$

$$\frac{\partial X_H}{\partial t} = e^{iHt/\hbar} \frac{\partial X(t)}{\partial t} e^{-iHt/\hbar}$$

$$= \frac{1}{i\hbar} \left([x(t) \cos \omega t, \frac{1}{2m} p(t)^2] \right.$$

$$\left. - \frac{1}{m\omega} [p(t) \sin \omega t, \frac{1}{2} m \omega^2 x(t)^2] \right)$$

$$+ x(t) (-\omega \sin \omega t) - \frac{1}{m} p(t) \cos \omega t$$

Since $[A, BC] = [A, B]C + B[A, C]$,

$$[x(t), p(t)^2] = [x(t), p(t)] p(t) + p(t) [x(t), p(t)]$$

$$= 2i\hbar p(t)$$

and $[p(t), x(t)^2] = [p(t), x(t)] x(t) + x(t) [p(t), x(t)]$

$$= -2i\hbar x(t)$$

so $\boxed{\frac{dX_H(t)}{dt}} = \frac{1}{i\hbar} \left(\frac{1}{2m} \cos \omega t \cdot 2i\hbar p(t) + \frac{1}{2} \omega \sin \omega t \cdot 2i\hbar x(t) \right)$

$$- \frac{1}{m} p(t) \cos \omega t - \omega x(t) \sin \omega t$$

$$= \boxed{0}$$

(c) $\boxed{\frac{dP_H(t)}{dt}} = \frac{1}{i\hbar} \left([p(t) \cos \omega t, \frac{1}{2} m \omega^2 x(t)^2] + [m\omega x(t) \sin \omega t, \frac{p(t)^2}{2m}] \right)$

$$- \omega p(t) \sin \omega t + m\omega^2 x(t) \cos \omega t$$

$$= \frac{1}{i\hbar} \left(\frac{1}{2} m \omega^2 \cos \omega t (-2i\hbar x(t)) + \frac{1}{2} \omega \sin \omega t (2i\hbar p(t)) \right)$$

$$- \omega p(t) \sin \omega t + m\omega^2 x(t) \cos \omega t$$

$$= \boxed{0}$$

$$\begin{aligned}
 4. (a) \frac{d^2}{dx^2} e^{-\alpha x^2} &= \frac{d}{dx} (-2\alpha x) e^{-\alpha x^2} = (-2\alpha + 4\alpha^2 x^2) e^{-\alpha x^2} \\
 \Rightarrow \boxed{E(\alpha)} &= \left(\frac{2\alpha}{\pi}\right)^{1/2} \int_{-\infty}^{\infty} dx \left[-\frac{\hbar^2}{2m} (-2\alpha + 4\alpha^2 x^2) + V(x) \right] e^{-\alpha x^2} \\
 &= \left(\frac{2\alpha}{\pi}\right)^{1/2} \left[\frac{\hbar^2}{2m} \left(2\alpha \int_{-\infty}^{\infty} dx e^{-\alpha x^2} - 4\alpha^2 \int_{-\infty}^{\infty} dx x^2 e^{-\alpha x^2} \right) + \int_{-\infty}^{\infty} dx V(x) e^{-\alpha x^2} \right] \\
 &= \boxed{\frac{\hbar^2}{2m} \underbrace{(2\alpha - \alpha)}_{=\alpha} + \left(\frac{2\alpha}{\pi}\right)^{1/2} \int_{-\infty}^{\infty} dx V(x) e^{-\alpha x^2}}
 \end{aligned}$$

$$\begin{aligned}
 (b) 0 &= \frac{dE(\alpha)}{d\alpha} \\
 &= \frac{\hbar^2}{2m} + \frac{1}{2} \alpha^{-1/2} \left(\frac{2}{\pi}\right)^{1/2} \int_{-\infty}^{\infty} dx V(x) e^{-2\alpha x^2} \\
 &\quad + \left(\frac{2\alpha}{\pi}\right)^{1/2} \int_{-\infty}^{\infty} dx V(x) e^{-2\alpha x^2} (-2x^2) \\
 \Rightarrow \frac{\hbar^2}{2m} &= \left(\frac{2\alpha}{\pi}\right)^{1/2} \cdot 2 \int_{-\infty}^{\infty} dx V(x) x^2 e^{-2\alpha x^2} \\
 &\quad - \left(\frac{1}{2\pi\alpha}\right)^{1/2} \int_{-\infty}^{\infty} dx V(x) e^{-2\alpha x^2} \\
 \Rightarrow \boxed{E(\alpha)} &= \left(\frac{2\alpha}{\pi}\right)^{1/2} \cdot 2\alpha \int_{-\infty}^{\infty} dx V(x) x^2 e^{-2\alpha x^2} \\
 &\quad - \frac{1}{2} \left(\frac{2\alpha}{\pi}\right)^{1/2} \int_{-\infty}^{\infty} dx V(x) e^{-2\alpha x^2} \\
 &\quad + \left(\frac{2\alpha}{\pi}\right)^{1/2} \int_{-\infty}^{\infty} dx V(x) e^{-2\alpha x^2} \\
 &= \boxed{\left(\frac{2\alpha}{\pi}\right)^{1/2} \int_{-\infty}^{\infty} dx V(x) \left(2\alpha x^2 + \frac{1}{2}\right) e^{-2\alpha x^2}}
 \end{aligned}$$

< 0 since $V(x) < 0$ in some region